

Discussion

The dimensions of the $\text{Se}(\text{SCN})_2$ molecule are $\text{Se-S} = 2.21 \text{ \AA.}$, $\text{S-C} = 1.69 \text{ \AA.}$, $\text{C-N} = 1.13 \text{ \AA.}$, $\angle \text{S-Se-S} = 101^\circ$, $\angle \text{Se-S-C} = 104^\circ$, dihedral angle for $\text{S-Se-S-C} = 79^\circ$. The sums of the covalent radii³ are $\text{Se-S} = 2.21 \text{ \AA.}$, $\text{S} = 1.81 \text{ \AA.}$, $\text{S-C} = 1.60 \text{ \AA.}$, $\text{C} \equiv \text{N} = 1.27 \text{ \AA.}$ and $\text{C} \equiv \text{N} = 1.15 \text{ \AA.}$ There is clearly no evidence for double-bond character in the Se-S bond. The S-C bond, however, is shorter than the normal single-bond distances found, for example, in dimethyl sulfide (1.82 \AA.),⁶ dimethyl sulfide (1.77 \AA.)⁷ and dimethyl trisulfide (1.78 \AA.)⁸ and corresponds to about 30% double-bond character. However, only small reliance can be placed on this result since the errors in the C coordinates could be as large as 0.15 \AA. The C-N distance given above is certainly too small, although one expects this distance to be less than the value of 1.22 \AA. found in HNCS .⁹ The mean distances found in the compound $[\text{Hg}(\text{SCN})_4][\text{Cu}(\text{en})_2]$ ¹⁰ are $\text{C-N} = 1.34 \text{ \AA.}$ and $\text{S-C} = 1.57 \text{ \AA.}$, which are in marked disagreement with our results. Since the SCN groups in the mercury compound appeared non-linear—a very unlikely result—we conclude that the interatomic distances are likely to be greatly in error.

The bond angles about Se and S are very close to those found in most sulfides, disulfides, trisulfides and comparable selenium compounds, although the

S-S-S angle in 2,2'-diiododiethyl trisulfide¹¹ appears to have the abnormally large value of 113° . The dihedral angle between the S-Se-S and Se-S-C planes, 79° , seems small. Pauling¹² predicts that this type of angle should be 90° or larger in sulfur compounds and has used this concept to explain the stability of S_8 . It is expected that selenium will show the same behavior, and Burbank¹³ found the average dihedral angle in α -selenium to be 102° . On the other hand, Donohue¹¹ found the S-S-S-C dihedral angle in 2,2'-diiododiethyl trisulfide to be 82° , although he did not place much reliance in this result. We do not believe that the errors in this determination are sufficient to cause a deviation of 11° or more in this angle. It is easily shown that $\partial\varphi/\partial x'_\text{C} = -15 \text{ deg./\AA.}$, $\partial\varphi/\partial y'_\text{C} = 21 \text{ deg./\AA.}$ and $\partial\varphi/\partial z'_\text{C} = 21 \text{ deg./\AA.}$, where φ is the dihedral angle and x'_C , y'_C , and z'_C are the coordinates of carbon in Å . The corresponding derivatives for nitrogen are even smaller. Hence, changes in the coordinates of both C and N of the order of 0.4 \AA. would be required to obtain a dihedral angle of 90° .

The packing of the molecules is such that the closest intermolecular distances are $\text{N} \cdots \text{Se}$ and $\text{N} \cdots \text{S}$. The shortest of these are 3.03 \AA. , 3.32 \AA. , 3.58 \AA. and 3.90 \AA. for $\text{N} \cdots \text{S}$, and 2.98 \AA. and 3.58 \AA. for $\text{N} \cdots \text{Se}$.

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Vapor Pressures of Inorganic Substances. XII. Zirconium Dioxide¹

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The vapor pressure of ZrO_2 has been determined by the method of Knudsen, between 2014 and 2290°K. ZrO_2 vaporizes undissociated as $\text{ZrO}_2(\text{g})$, the heat of vaporization being $\Delta H_{298}^\circ = 153.6 \pm 1 \text{ kcal}$. The vapor pressure can be represented by the equation $\log p(\text{atm.}) = -(34,383/T) - 7.98 \times 10^{-4}T + 11.98$. The heat of dissociation of $\text{ZrO}_2(\text{g})$ into gaseous atoms was found to be $D_0 = 365 \pm 5 \text{ kcal}$.

Introduction

No experimental determination of the vapor pressure of zirconium dioxide is reported in the literature, and the question concerning which gaseous species are present when ZrO_2 vaporizes is still not settled.

The only available information on the vapor pressure of ZrO_2 is that of Searcy² who estimated a vapor pressure of $1.5 \times 10^{-7} \text{ atm.}$ at 2073°K. from measurements of Zintl, Morawietz and Gastingner.³ Mott⁴ reported the boiling point at 4570°K. Starodubtsev⁵ observed ZrO_2^+ in a mass spectro-

graph, which indicates the existence of the gaseous ZrO_2 molecule.

In the present investigation, the Knudsen effusion method was used to determine the pressure over pure ZrO_2 and over a mixture of ZrO_2 and Zr.

Apparatus and Experimental Procedure

Two tantalum Knudsen cells were used, of dimensions 1 in. diameter and $3/4$ in. high, with an orifice diameter of $1/16$ in. The area of the orifice was corrected for thermal expansion by using the thermal expansion coefficient of tantalum determined previously in this Laboratory.⁶ ZrO_2 was introduced into one of the cells, while a mixture of ZrO_2 and Zr was introduced into the other. Both cells were degassed for two hours at 2000° before the runs were made.

The samples were heated in our metal cell,⁷ and the temperature was measured with a disappearing-filament optical pyrometer, calibrated against a standard tungsten-ribbon

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amp. The standard lamp had been calibrated by the National Bureau of Standards, and in the range of the present investigation was reported to have a maximum uncertainty of 5°.

For the small temperature fluctuation during each run, and for evaporation during the initial heating and cooling periods, the averaging method^{8,9} was used. To correct for the thickness of the effusion hole (0.01 in.), all measured weight losses were multiplied by the factor $1/0.89 = 1.1236$.^{10,11}

The ZrO₂ was obtained from the Oak Ridge National Laboratories, and its hafnium content was less than 0.025%. The Zr was in powder form; it was obtained from Charles Hardy Inc., New York, and its Zr metal content was better than 99.5%.

Experimental Results and Discussion of the Data

The experimental data from pure ZrO₂ are presented in Table I, and those for the stoichiometric ZrO₂ + Zr mixture are given in Table II. The data from the two tables are plotted in Fig. 1. The pressure was calculated from the rate of effusion data, using the equation $p = m \sqrt{2\pi RT/M}$, where p is the pressure in atmospheres, R is the molar gas constant, T is the absolute temperature, m is the rate of effusion in g. cm.²/sec., and M is the molecular weight of the vapor. The molecular species was taken as ZrO₂, according to the following discussion.

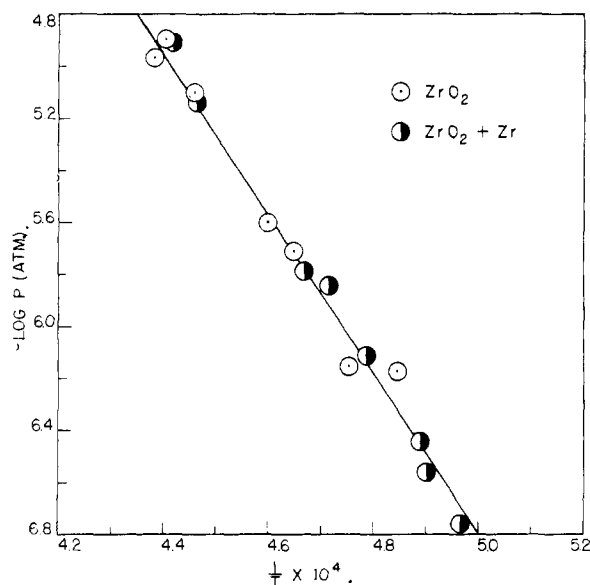


Fig. 1.—Vapor pressure of ZrO₂ and ZrO₂ + Zr.

The rates of evaporation were found to be the same for both cases, indicating that ZrO₂ vaporizes as ZrO₂. If dissociation into ZrO and O occurred, then the pressure above the ZrO₂ + Zr mixture would be much larger than that above the pure ZrO₂. Vaporization into Zr metal and oxygen is

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TABLE I

VAPOR PRESSURE OF ZrO₂

Run	Temp., °K.	Eff. time, sec.	Area, cm. ²	Wt. loss, g.	Evapn. rate, g./cm. ² /sec. × 10 ⁵	p , atm. × 10 ⁶	$-\log p$
1	2064	47,903	0.0210	0.0073	0.7257	0.670	6.174
3	2104	44,971	.0210	.0081	0.8577	0.799	6.156
2	2151	34,206	.0210	.0148	2.0603	1.946	5.711
5	2175	32,658	.0210	.0181	2.6392	2.500	5.602
4	2244	24,071	.0211	.0418	8.2300	7.922	5.101
7	2271	17,435	.0211	.0481	13.075	12.660	4.898
6	2282	16,304	.0211	.0381	11.075	10.750	4.969

TABLE II

VAPOR PRESSURE OF ZrO₂ + Zr

Run	Temp., °K.	Eff. time, sec.	Area, cm. ²	Wt. loss, g.	Evapn. rate, g./cm. ² /sec. × 10 ⁵	p , atm. × 10 ⁶	$-\log p$
3	2014	28,851	0.02041	0.0011	0.1868	0.174	6.759
1	2040	57,022	.02042	.0035	.3006	.275	6.561
4	2046	70,123	.02042	.0056	.3911	.361	6.443
5	2089	43,823	.02043	.0074	.8265	.769	6.114
2	2121	38,680	.02044	.0121	1.5304	1.436	5.843
6	2142	50,873	.02045	.0184	1.7686	1.666	5.778
7	2240	21,963	.02049	.0339	7.5330	7.253	5.139
8	2290	17,811	.02051	.0463	12.6744	12.330	4.909

excluded because the measured rates of evaporation are 100 times larger than those for zirconium.⁹ The dissociation into gaseous oxygen and solid Zr is excluded by the following facts: (1) Twice as much material was vaporized from each cell as the amount of oxygen contained in the ZrO₂ and in the ZrO₂ + Zr mixture, and (2) X-ray patterns taken after the runs showed only ZrO₂ lines in material from the ZrO₂ cell, and expanded Zr lines (saturated with oxygen) in the material from the ZrO₂ + Zr cell.

The data are plotted in Fig. 1.

To obtain ΔH° and the equation for the vapor pressure, our data were treated by a sigma plot [$\Delta F^\circ/T + \Delta \alpha \ln T = \Sigma = \Delta H^\circ/T + 1$]. The necessary specific heat data were obtained from Kelley,¹² who gives $C_p = 17.80 - 4.00 \times 10^{-5} \times T^{-2}$ for solid ZrO₂ and we estimated the specific heat of gaseous ZrO₂ to be 10 cal. mole⁻¹ degree⁻¹. From the sigma plot we obtain $\Delta H_0^\circ = 157.3 \pm$ kcal. and

$$\Delta H = 157,300 - 7.80T - 4 \times 10^5 T^{-1}$$

After expanding the term $-7.80 \ln T + 2 \times 10^5 T^{-2}$ as a linear function of the temperature we obtain

$$\log p(\text{atm.}) = -\frac{34,383}{T} - (7.98 \times 10^{-4})T + 11.98$$

By using the value determined in the Laboratory⁹ for the heat of evaporation of Zr, and the heat of formation of solid ZrO₂ given by Humphrey¹³ a value of 368 ± 5 kcal./mole was obtained for the heat of dissociation of gaseous ZrO₂ into gaseous atoms of Zr and O.

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